

## Explanation of the anomalies of water

### 1 High melting point (0°C, c.f. CHCl<sub>3</sub> -63°C)

In ice (1h), all water molecules participate in four hydrogen bonds (two as donor and two as acceptor) and are held relatively static. In liquid water, some hydrogen bonds must be broken to allow the molecules to move around. The large energy required for breaking these bonds must be supplied during the melting process; only a relatively minor amount of energy is reclaimed from the change in volume. The free energy change ( $\Delta G = \Delta H - T\Delta S$ ) must be zero at the melting point. As temperature is increased, the amount of hydrogen bonding in liquid water decreases and its entropy increases. Melting will only occur when there is sufficient entropy change to provide the energy required for the bond breaking. The low entropy (high organization) of liquid water causes this melting point to be high. [Back]

### 2 High boiling point (100°C, c.f. CHCl<sub>3</sub> 61°C)

There is considerable hydrogen bonding in liquid water which prevents water molecules from being easily released from the water's surface. This reduces the vapor pressure. Boiling cannot occur until this vapor pressure equals the external pressure. This occurs at a consequentially higher temperature. [Back]

### 3 High critical point (374°C, c.f. CH<sub>3</sub>CH<sub>3</sub> 32°C)

The critical point can only be reached when the interactions between the water molecules fall below a certain threshold level. Due to the strength and extent of the hydrogen bonding, much energy is needed to cause a reduction in molecular interaction and this requires higher temperatures. Even close to the critical point, a considerable number of hydrogen bonds remain, albeit bent, elongated and no longer tetrahedrally arranged [92]. [Back]

### 4 High surface tension (72.75 mJ/m<sup>2</sup>, c.f. CCl<sub>4</sub> 26.6 mJ/m<sup>2</sup> at 20°C)

Water molecules at the liquid-gas surface have lost potential hydrogen bonds directed at the gas phase and are pulled towards the underlying bulk liquid water by the remaining stronger hydrogen bonds [214]. Energy is required to increase the surface area, so it is minimized and held under tension. As the forces between the water molecules are several and relatively large on a per-mass basis, compared to those between most other molecules, the surface tension is large. Lowering the temperature greatly increases the hydrogen bonding, causing increased surface tension. If a small drop of water (typically 1 mm diameter) is coated in a fine (typically 20 nm diameter) hydrophobic dust then the drop can roll and bounce without leakage [225], and aqueous spheres can even float on water. Capillarity holds the dust at the air-liquid interface with the effect being due to the high surface tension. The interesting absence of any apparent anomaly in the surface tension/temperature behavior (although unexpectedly, there is no deviation associated with the density anomaly at 4°C, but there is an inflection point at about 250°C [427]) is due to surface enrichment with clusters that are different from those present in the bulk. These interfacial water molecules have recently been found experimentally to possess a 6% expanded but weaker hydrogen bonded structure [415], due perhaps to less hydrogen bond polarization, itself caused by the missing hydrogen bonds. The affinity of chaotropic

for this expanded water structure may help explain their shallow minima in the surface tension at very low ionic concentrations (*i.e.* the Jones-Ray effect); higher concentrations destroying the ordered structuring required. The greater than expected drop in surface tension with temperature increase has been quantitatively explained using spherically symmetrical water clustering [376].

It is interesting to note that surfactants lower the surface tension because they prefer to sit in the surface, attracting the surface water molecules in competition to the bulk water hydrogen bonding and so reducing net forces away from the surface (*i.e.* the surface tension). In contrast, most cations and anions prefer to be fully hydrated in the bulk liquid water so adding to the attractive forces on the surface water molecules, consequently increasing the surface tension. [Back]

### **5 High viscosity (0.89 cP, c.f. pentane 0.22 cP, at 25°C)**

The viscosity of a liquid is determined by the ease with which molecules can move relative to each other, depends on the forces holding the molecules together. This cohesivity is large in water due to its extensive three-dimensional hydrogen bonding. It should be noted that although the viscosity is high, it is not so high that it causes too much difficulty being moved around within organisms. The Arrhenius energy of activation for viscous flow is similar to the hydrogen bond energy ( $\text{H}_2\text{O}$ ,  $21.5 \text{ kJ mol}^{-1}$ ;  $\text{D}_2\text{O}$ ,  $24.7 \text{ kJ mol}^{-1}$ ;  $\text{T}_2\text{O}$ ,  $26.5 \text{ kJ mol}^{-1}$ , all calculated from [73]; all at  $0^\circ\text{C}$  and all more than doubling at  $-30^\circ\text{C}$ ) [Back]

### **6 High heat of vaporization ( $40.7 \text{ kJ mol}^{-1}$ , c.f. $\text{H}_2\text{S}$ $18.7 \text{ kJ mol}^{-1}$ )**

There is still considerable hydrogen bonding ( $\sim 75\%$ ) in water at  $100^\circ\text{C}$ . As effectively all these bonds need to be broken (very few indeed remaining in the gas phase), there is a great deal of energy required to convert water to gas, where the water molecules are effectively separated. The increased hydrogen bonding at low temperatures causes higher heats of vaporization (e.g.  $44.8 \text{ kJ mol}^{-1}$ , at  $0^\circ\text{C}$ ). The high heat of vaporization also causes water to have an anomalously low ebullioscopic constant (*i.e.* effect of solute on boiling point elevation,  $0.51 \text{ K kg/mol}$ , c.f.  $\text{CCl}_4$   $4.95 \text{ K kg/mol}$ ). Water also has an anomalously high entropy of vaporization ( $109 \text{ J K}^{-1} \text{ mol}^{-1}$ , cf. Trouton's constant  $85 \text{ J K}^{-1} \text{ mol}^{-1}$ ) due to the additional hydrogen-bond order lost on vaporization. As both values are anomalously high, their ratio ( $\text{DH}_{\text{vap}}/\text{DS}_{\text{vap}}$ ) is not anomalous. The high heat of vaporization also gives rise to an anomalously high heat of sublimation. [Back]

### **7 Shrinks on melting**

When water freezes at  $0^\circ\text{C}$ , at atmospheric pressure, its volume increases by about 9%. If the melting point is lowered by increased pressure, the increase in volume on freezing is even greater (e.g. 13% at  $-20^\circ\text{C}$ ). The structure of ice (Ih) is open with a low packing efficiency where all the water molecules are involved in four straight tetrahedrally-oriented hydrogen bonds; for comparison, solid hydrogen sulfide has a face centered cubic closed packed structure with each molecule having twelve nearest neighbors [119]. On melting, some of these ice (Ih) bonds break, others bend and the structure undergoes a partial collapse. This is different from what happens with most solids (but similar to other tetrahedrally arranged solids such as silicon and silica) where the extra movement available in the liquid phase requires more space and therefore melting is accompanied by expansion. If the increase in volume on freezing is prevented, an increased pressure of up to 25 MPa may be generated; easily capable of bursting pipes in Winter<sup>4</sup>. Melting ice, within a fixed and sealed volume, may result in an apparently superheated state where the metastable iso-dense liquid water is stretched, relative to its equilibrium state at the effectively-negative pressure, due to its cohesiveness.

Consequently, the  $ES \rightleftharpoons CS$  equilibrium is shifted towards the more-open ES structure.

In contrast, it should be noted that the high pressure ices (ice III, ice V, ice VI and ice VII) all expand on melting to form liquid water (under high pressure). [Back]

### **8 High density that increases on heating (up to 3.984°C); the density anomaly**

The high density of liquid water is due mainly to the cohesive nature of the hydrogen-bonded network. It reduces the free volume and ensures a relatively high-density, compensating for the partial open nature of the hydrogen-bonded network. Its density, however, is not as great as that of closely packed, isoelectronic, liquid neon ( $1207 \text{ kg m}^{-3}$  at 27 K). The anomalous temperature-density behavior of water can be explained as previously [13, 14] utilizing the range of environments within whole or partially formed clusters with different degrees of dodecahedral puckering. The density maximum is brought about by the opposing effects of increasing temperature, causing both structural collapse that increases density and thermal expansion that lowers density. At lower temperatures there is a higher concentration of ES whereas at higher temperatures there is more CS and fragments, but the volume they occupy expands with temperature. The change from ES to CS as the temperature rises is accompanied by positive changes in entropy and enthalpy due to the less ordered structure and greater hydrogen bond bending respectively.

A similar but unrelated phenomena occurs in hexagonal, cubic and amorphous ices which all expand slightly with cooling below about 50 K. This appears to be due to alteration in the net bending motion of three tetrahedral hydrogen bonded molecules with temperature, as higher frequency modes are reduced [209].

A further strange anomaly in water is that at a temperature below the maximum density anomaly there is a minimum density anomaly so long as no phase change occurs. This is expected to lie just below the minimum temperature accessible on supercooling (232 K, [215]) and close to where both maximum ES structuring and compressibility occur, with the liquid density close to that of hexagonal ice. [Back]

### **9 The number of nearest neighbors increases on melting**

Each water molecule in hexagonal ice has four nearest neighbors. On melting, the partial collapse of the hydrogen bonded network allows unbonded molecules to approach more closely so increasing this number. Normally in a liquid the movement of molecules, and the extra space they find themselves in, means that it becomes less likely that they will be found close to each other; e.g. argon has exactly twelve nearest neighbors in the solid state but only an average of about ten on melting. [Back]

### **10 Nearest neighbors increase with temperature**

If a water molecule is in a fully hydrogen-bonded structure with strong and straight hydrogen bonds (such as hexagonal ice) then it will only have four nearest neighbors. In the liquid phase, molecules approach more closely due to the partial collapse of the open hydrogen bonded network. As the temperature of liquid water increases, the continuing collapse of the hydrogen bonded network allows unbonded molecules to approach more closely so increasing the number of nearest neighbors. This is in contrast to normal liquids where the increasing kinetic energy of molecules and space available due to expansion, as the temperature is raised, means that it becomes less likely that molecules will be found close to each other. [Back]

### **11 Pressure reduces its melting point (13.35 MPa gives a melting point of -1°C)**

Increasing pressure normally promotes liquid freezing, shifting the melting point to higher temperatures. This is shown by the forward sloping liquid/solid line in the phase diagram. In water, this line is backward sloping. As the pressure increases, the water equilibrium shifts towards a collapsed structure (e.g. CS) with higher entropy. This lowers the melting free energy change ( $\Delta G = \Delta H - T\Delta S$ ) such that it will be zero (i.e. at the melting point) at a lower temperature. The minimum temperature that liquid water can exist without ever freezing is  $-21.985^\circ\text{C}$  at 209.9 MPa; at higher pressures water freezes at higher temperatures. Stretching has the reverse effect; ice melting at  $+6.5^\circ\text{C}$  at about -95 MPa negative pressure within stretched microscopic aqueous pockets in mineral fluorite [243].

It should be noted that ice skating does not produce sufficient pressure to lower the melting point significantly; the slipperiness being generated by frictional heating. [Back]

## 12 Pressure reduces the temperature of maximum density

Increasing pressure shifts the water equilibrium towards a more collapsed structure (e.g. CS). So although pressure will increase the density of water at all temperatures (flattening the temperature density curve), it has a disproportionate effect at lower temperatures. The result is a shift in the temperature of maximum density to lower temperatures. At high enough pressures the density maximum is shifted to below  $0^\circ\text{C}$  ( $\sim$  200 MPa) and, at above 200 MPa, it cannot be observed. A similar effect may be caused by increasing salt concentration, which behaves like increased pressure in breaking up the low-density clusters. Thus in 0.3 Molal NaCl the temperature of freezing and maximum density coincide at  $-1.33^\circ\text{C}$ . Higher salt concentrations reduce the temperature of maximum density such that it is only accessible in the supercooled liquid. The stronger hydrogen bonding in  $\text{D}_2\text{O}$  gives rise to a smaller shift in the temperature of maximum density with respect to increasing pressure.

At very high negative pressures (i.e. increased stretching of liquid water), the temperature of maximum density may be reduced as the hydrogen bonds are stretched to breaking point; the temperature of maximum density showing a maximum with respect to pressure in the negative pressure region [419]. [Back]

## 13 $\text{D}_2\text{O}$ and $\text{T}_2\text{O}$ differ significantly from $\text{H}_2\text{O}$ in their physical properties

These heavier forms of water hydrogen bond more strongly than normal water. This causes many of their properties (such as the viscosity, self-diffusion coefficient, protein solubility and toxicity [424]) to be different from those expected from a simple consideration of their increased mass (e.g. the  $\text{D}_2\text{O}/\text{H}_2\text{O}$  viscosity ratio rises from about 1.16 at  $100^\circ\text{C}$  to around 2.0 in deeply supercooled water [23b]). This difference appears as a shift in the equilibrium position equivalent to a slight increase in temperature [425]; e.g. viscosity data have been reconciled if the temperatures are shifted by  $6.498^\circ\text{C}$  and  $8.766^\circ\text{C}$  for  $\text{D}_2\text{O}$  and  $\text{T}_2\text{O}$  respectively [7].

Mixtures of  $\text{H}_2\text{O}$  and  $\text{D}_2\text{O}$  equilibrate to form HDO:



which is close to a total randomisation of the hydrogen atoms (i.e. equal concentrations of HOH, HOD, and DOD giving  $K_{\text{eq}} = 4$ ). Even the properties of HDO deviate from those expected from a consideration of the properties of  $\text{H}_2\text{O}$  and  $\text{D}_2\text{O}$  [126]. [Back]

## 14 Large viscosity increase as the temperature is lowered

The water cluster equilibrium shifts towards the more open structure (e.g. ES) as the temperature is lowered.

This structure is formed by stronger hydrogen bonding. In turn, this creates larger clusters and reduces the ease of movement (increasing viscosity). Self-diffusion is also greatly reduced at lower temperatures which anomalously decreases as the density decreases (see below). [Back]

### **15 Viscosity decreases with pressure (at temperatures below 33°C)**

Viscous flow occurs by molecules moving through the voids that exist between them. As the pressure increases, the volume decreases and the volume of these voids reduces, so normally increasing the viscosity. Water's pressure-viscosity behavior can be explained by the increased pressure (up to about 150 MPa) causing deformation, so reducing the strength of the hydrogen-bonded network, which is also partially responsible for the viscosity. This reduction in cohesivity more than compensates for the reduced void volume. At higher pressures water behaves as a normal liquid; viscosity increasing with pressure. Self-diffusion is also affected by pressure where (at low temperatures) both the translational and rotational motion of water anomalously increase as the pressure increases (see below). [Back]

### **16 Low compressibility (0.46 GPa<sup>-1</sup>, c.f. CCl<sub>4</sub> 1.05 GPa<sup>-1</sup>, at 25°C)**

It may be thought that water should have a high compressibility ( $k_T = -[\partial V / \partial P]_T / V$ ) as the large cavities in liquid water allows plenty of scope for the water structure to collapse under pressure without water molecules approaching close enough to repel each other. The deformation causes the growth in the radial distribution function peak at about 3.5 Å with increasing pressure [51] (and temperature [50]), due to the collapse in structure. The low compressibility of water is due to water's high-density, again due to the cohesive nature of the extensive hydrogen bonding. This reduces the free space (compared with other liquids) to a greater extent than the contained cavities increase it. At low temperatures D<sub>2</sub>O has a higher compressibility than H<sub>2</sub>O (4% higher at 10°C but only 2% higher at 40°C [188]) due to its stronger hydrogen bonding producing an ES ⇌ CS equilibrium shifted towards the more-open ES structure. [Back]

### **17 Compressibility drops as temperature increases (up to a minimum at about 46.5°C)**

In a typical liquid the compressibility decreases as the structure becomes more compact due to lowered temperature. In water, the cluster equilibrium shifts towards the more open structure (e.g. ES) as the temperature is reduced due to it favoring the more ordered structure (DG becomes more positive). As the water structure is more open at these lower temperatures, the capacity for it to be compressed increases. This effect is not a simple dependency on density, however, or else the minimum at 46.5°C for isothermal (i.e. without change in temperature) compressibility ( $k_T = -[\partial V / \partial P]_T / V$ ) and the minimum at 64°C for adiabatic (i.e. without loss or gain of heat energy, also isentropic) compressibility ( $k_S = -[\partial V / \partial P]_S / V$  [112]) would be at the density minimum (4°C). Compressibility depends on fluctuations in the specific volume and the will be large where water molecules fluctuate between being associated, or not, with a more open structure and between the different environments within the water clusters. At sufficiently low temperature, there will be a maximum in this compressibility-temperature relationship, so long as no phase change occurs. This is expected to lie just below the minimum temperature accessible on supercooling (232 K, [215]) close to the temperature of minimum density. [Back]

### **18 Water has a low thermal expansivity (0.00021/°C, c.f. CCl<sub>4</sub> 0.00124/°C at 20°C)**

As the temperature increases, the cluster equilibrium shifts towards the more collapsed structure (e.g. CS

which reduces any increase in volume due to the increased kinetic energy of the molecules. Normally the higher the volume a molecule occupies, the larger is the disorder (entropy). Thermal expansivity ( $\alpha_p = \left(\frac{\partial V}{\partial T}\right)_p/V$ ) depends on the product of the fluctuations in these factors. In water, however, the more open structure (e.g. ES) is also more ordered (i.e. has lower entropy). [Back]

### 19 Water's thermal expansivity reduces increasingly (becoming negative) at low temperatures.

Supercooled and cold ( $< 3.984^\circ\text{C}$ ) liquid water contracts on heating. As the temperature decreases, the equilibrium shifts towards the expanded, more open, structure (e.g. ES), which more than compensates for the decrease in volume due to the reduction in the kinetic energy of the molecules. It should be noted that this behavior requires that the thermodynamic work ( $dW$ ) equals  $-p dV$  rather than the usual  $+p dV$  (pressure change in volume) [404]. [Back]

### 20 Speed of sound increases with temperature (up to a maximum at $73^\circ\text{C}$ )

Sound is a longitudinal pressure wave, whereby the energy is propagated as deformations in the media but molecules then return to their original positions and are not propagated. The propagation of a sound wave depends on the transfer of vibration from one molecule to another. The speed ( $u = \left(\frac{\partial r}{\partial P}\right)_S$ ) is inversely proportional to the square root of the adiabatic compressibility times density term. The anomalous nature of both these physical properties are described above (compressibility, density). At low temperatures both compressibility and density are high, so causing a lower speed of sound. As the temperature increases the compressibility drops and goes through a minimum whereas the density goes through a maximum and then drops. Combination of these two properties leads to the maximum in the speed of sound. Increasing the pressure increases the speed of sound and shifts the maximum to higher temperatures, both in line with the effect on the density.

Water has a second anomaly (called 'fast sound') concerning the speed of sound. Over a range of high frequencies liquid water behaves as though frozen and sound travels at about twice its normal speed ( $\sim 32 \text{ s}^{-1}$ ). [Back]

### 21 High specific heat capacity ( $C_V$ and $C_P$ , $4.18 \text{ J g}^{-1} \text{ K}^{-1}$ ) c.f. pentane $1.66 \text{ J g}^{-1} \text{ K}^{-1}$ , at $25^\circ\text{C}$ )

As water is heated, the increased movement of water causes the hydrogen bonds to bend and break. As the energy absorbed in these processes is not available to increase the kinetic energy of the water, it takes considerable heat to raise water's temperature. Also, as water is a light molecule there are more molecules per gram, than most similar molecules, to absorb this energy. Heat absorbed is given out on cooling, so allowing water to act as a heat reservoir, buffering against changes in temperature. [Back]

### 22 Water has over twice the specific heat capacity of ice or steam (c.f. benzene where $C_{P \text{ liquid}} = 1 C_{P \text{ solid}}$ )

The specific heats of polar molecules do increase considerably on melting but water shows a particularly large increase. As water is heated, much of the energy is used to bend the hydrogen bonds; a factor not available in the solid or gaseous phase. This extra energy causes the specific heat to be greater in liquid water. The presence of this large specific heat offers strong support for the extensive nature of the hydrogen-bonded network of liquid water. [Back]

### 23 Specific heat capacity ( $C_p$ ) has a minimum (36°C)

The specific heat capacity has a shallow minimum at about 36°C with a particularly steep negative slope below 273°C [15]. The water cluster equilibrium shifts towards less structure (e.g. CS) and higher enthalpy as the temperature is raised.  $C_p$  is the heat capacity at constant pressure defined by  $C_p = (\partial H / \partial T)_p$  (change in enthalpy with temperature). The extra positive  $\partial H$  due to the shift in equilibrium (at low temperatures) as temperature is raised causes a higher  $C_p$  than otherwise, particularly at supercooling temperatures where a much larger shift occurs. This addition to the  $C_p$ , as the temperature is lowered, is greater than the 'natural' fall expected, so causing a minimum to be created. Note that  $C_v$  (the heat capacity at constant volume,  $C_v = (\partial U / \partial T)_v$ ) shows no such anomaly, as most of the anomalous enthalpy change is associated with the anomalous volume change. It is expected that the large specific heat changes with temperature at low temperatures will be reduced at higher pressures and the specific heat minimum will shift to lower temperatures. At sufficiently low temperature, there must be a maximum in this compressibility-temperature relationship, so long as no phase change occurs. This is expected to lie just below the minimum temperature accessible on supercooling (232 K, [215]). [Back]

### 24 NMR spin-lattice relaxation times are very small at low temperatures

NMR spin-lattice relaxation time depends on the degree of structure. As the water cluster equilibrium shifts towards a stiffer, tetrahedrally organized, structure (e.g. ES) as the temperature is lowered, the NMR spin-lattice relaxation time reduces far more than would otherwise be expected. This effect can be partially reversed by increasing the pressure, which reduces the degree of structure. [Back]

### 25 Solutes have varying effects on properties such as density and viscosity

Solutes will interfere with the cluster equilibrium by favoring either open or collapsed structures. Any effect will cause the physical properties of the solution, such as density or viscosity, to change. Solutes have a larger than expected effect on both the cryoscopic (i.e. effect of solute on freezing point depression, 1.86 K kg mol<sup>-1</sup> c.f. CCl<sub>4</sub> 30 K kg mol<sup>-1</sup>) and ebullioscopic constants. [Back]

### 26 Solutions are not ideal

Ideality depends on the structure of the solvent being unaffected by the solute. Water is not even close to being a homogeneous phase on the molecular level. Local clustering will be effected by the presence of solutes, so changing the nature of the water. Even solutions of HDO in H<sub>2</sub>O do not behave ideally. [Back]

### 27 X-ray diffraction shows an unusually detailed structure

This is simply explained by the presence of ordered clustering within the liquid phase. [Back]

### 28 Supercooled water has two phases and a second critical point at about -50°C

As water is supercooled it converts mainly into its expanded form (e.g. ES) at ambient pressures, which at low enough temperatures (< -38°C) may result in it forming metastable low-density amorphous ice (LDA)

although normally it will form hexagonal ice at this temperature). If the pressure on LDA is increased above about 200 MPa then LDA undergoes a 30% collapse forming metastable high-density amorphous ice (HDA) but notably in a continuous process without breaking the hydrogen bonds [394]. This phase change can continue to higher temperatures (so creating a second critical point, [45]) as neither of these phases is stable in the presence of liquid water although they may convert into their metastable supercooled liquid forms. [E]

## **29 Liquid water is easily supercooled**

It may be expected that the directional hydrogen bonding capacity of water would reduce its tendency to supercool as it would encourage the regular structuring in cold liquid. Liquid water, however, is easily supercooled down to about  $-25^{\circ}\text{C}$  and with more difficulty down to about  $-38^{\circ}\text{C}$  with further supercooling possible, in tiny droplets, down to about  $-70^{\circ}\text{C}$ . Water, supercooled down to  $-37.5^{\circ}\text{C}$ , is sustained in storm clouds and the condensed clouds formed by aircraft. Rather strangely, at the limit of this supercooling (also known as the homogeneous freezing point) the water activity is always 0.305 lower than that of water measured at the same temperature [457].

Liquid water may be supercooled to about  $-92^{\circ}\text{C}$  at 210 MPa. Deeply supercooled liquid water can be produced from glassy amorphous ice between  $-123^{\circ}\text{C}$  and  $-149^{\circ}\text{C}$  [74] and may coexist with cubic ice up to  $-63^{\circ}\text{C}$  [137]. This behavior is particularly anomalous as the liquid (deeply supercooled water) is a 'strong' liquid (compared with supercooled water that is a 'fragile' liquid [493]) that changes to crystalline solid (ice) on increasing the temperature whilst keeping the pressure constant.

As water is cooled, the cluster equilibrium shifts towards the more open structure (e.g. ES) with higher viscosity. In order for crystallization to occur at least 3 - 4 unit cells worth of water molecules have to come together in the correct orientation.<sup>1</sup> The formation of icosahedral clusters interferes with this process whilst not allowing cluster crystallization due to their five fold symmetry. Lowering the temperature further, which should encourage crystallization, is partially counteracted by the increase in icosahedral clustering. The presence of ES clusters is, in principle, in agreement with computer simulation studies requiring the presence of metastable states [216]. Methods that break the hydrogen bonding in these clusters, such as ultrasonics [296], cause the supercooled water to freeze.

A possible explanation of the existence of low-temperature-range supercooled water (124-150 K) may be the formation of strands of icosahedral structures. This model can also explain the high viscosity and strong (low specific heat) liquid behavior of this extremely supercooled water [215]. [Back]

## **30 Solid water exists in a wide variety of stable (and metastable) crystal and amorphous structures**

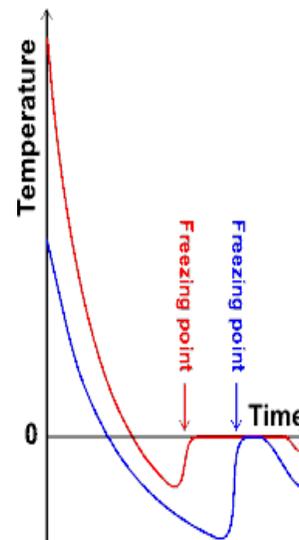
The ability for water to form extensive networks of hydrogen bonds increases the number of solid phases possible. The open structure of hexagonal ice ( $19.65\text{ cm}^3/\text{mol}$ ), which contains only about  $7.5\text{ cm}^3/\text{mol}$  of water molecules, gives plenty of scope for different arrangements of the water molecules as the structure is compressed. For comparison, hydrogen sulfide has only four distinct solid phases [119]. [Back]

## **31 Hot water may freeze faster than cold water; the Mpemba effect**

The ability of hot water to freeze faster than cold seems counter-intuitive as it would seem that hot water must first become cold water and therefore the time required for this will always delay its freezing relative to cold water. However experiments show that hot water (e.g.  $90^{\circ}\text{C}$ ) does often appear to freeze faster than the same amount of cold water (e.g.  $18^{\circ}\text{C}$ ) under otherwise identical conditions [158]. This has been recognised even

far back as Aristotle in the 4<sup>th</sup> century BC but was brought to the attention of the scientific community by perseverance of a Tanzanian schoolboy (Erasto Mpemba) who refused to reject his own evidence, or bow to disbelieving mockery, that he could freeze ice cream faster if he warmed it first.

A number of explanations have been put forward but the most likely scenario (described in [158]) is that the degree of supercooling is greater in initially-cold water than initially-hot water. The initially-hot water appears to freeze at a higher temperature (less supercooling) but less of the apparently frozen ice is solid and a considerable amount is trapped liquid water. Initially-cold water freezes at a lower temperature to a more completely solid ice with less included liquid water; the lower temperature causing intensive nucleation and a faster crystal growth rate. If the freezing temperature is kept about  $-6^{\circ}\text{C}$  then the initially-hot water is most likely to (apparently) freeze first. If freezing is continued, initially-cold water always completely freezes before initially-hot water. Why initially-cold water supercools more is explained in terms of the clustering of water. Icosahedral clusters do not readily allow the necessary arrangement of water molecules to enable hexagonal ice crystal initiation; such clustering is the cause of the facile supercooling of water. Water that is initially-cold will have the maximum (equilibrium) concentration of such icosahedral clustering. Initially-hot water has lost much of its ordered clustering and, if the cooling time is sufficiently short, this will not be fully re-attained before freezing.



Experiments on low-density water around macromolecules has shown that such clustering processes may some time [4]. It is also possible that dissolved gases may encourage supercooling by (1) increasing the  $\rho$  of structuring, by hydrophobic hydration, in the previously-cold water relative to the gas-reduced previously-hot water (the critical effect of low concentrations of dissolved gas on water structure is reported [294]; re-equilibration taking several days) and (2) increasing the pressure as gas comes out of solution when the water starts to crystallize, so lowering the melting point and reducing the tendency to freeze (see guestbook). Also, the presence of tiny gas bubbles (produced on heating) may increase the rate of nucleation so reducing supercooling [428]. [Back]

### 32 The refractive index of water has a maximum value at just below $0^{\circ}\text{C}$ .

The refractive index of water ( $\lambda = 589.26 \text{ nm}$ ) rises from an estimated 1.33026 at  $-30^{\circ}\text{C}$  to a maximum value just below  $0^{\circ}\text{C}$  (1.33434) before falling ever increasingly to 1.31854 at  $100^{\circ}\text{C}$  [310]. This may be explained by the mixture model [60] applied to the change from ES to CS as the temperature rises; ES possessing a lower refractive index than CS. Most of the effect is due to the density difference between ES and CS. High density produces higher refractive index such that the refractive index temperature maximum lies close to density maximum, with the small difference due to the slightly different effect of temperature on the specific refractions of ES and CS. Although not considered anomalous, it is interesting to note that ice has the low refractive index (1.31,  $\lambda = 589 \text{ nm}$ ) of any known crystal. [Back]

### 33 The solubilities of non-polar gases in water decrease with increasing temperature to a minimum and then rise.<sup>2</sup>

Non-polar gases are poorly soluble in water. Their hydration may be considered as the sum of two processes

(A) the endothermic opening of a clathrate pocket in the water, and (B) the exothermic placement of a molecule in that pocket, due to the multiple van der Waals interactions. In water at low temperatures, the energy required by process (A) is very small as such pockets may be easily formed within the water clust (by CS @ ES)<sup>3</sup>. The solubilization process is therefore exothermic and (as predicted by Le Chatelier's principle) solubility decreases with temperature rise. At high temperatures (often requiring high pressure), natural clustering is much reduced causing greater energy to be required for opening of the pocket in the water. The solubilization process therefore becomes endothermic and (as predicted by Le Chatelier's principle) solubility goes through a minimum before increasing with temperature rise. The more attractive solute-water van der Waals interactions (due both to atomic number dependency and goodness of fit with the clathrate pocket), the greater the inherent exothermic nature of the process and therefore the higher the temperature minimum (e.g. He, 32°C; Ne, 50°C; Ar, 96°C; Kr, 101°C; Xe, 123°C) and the greater the temperature range of negative temperature solubility coefficient. Similarly Henry's constants (= partial pressure/mole fraction) exhibit decreasing maximas at 30°C, 50°C, 90°C, 108°C and 110°C in H<sub>2</sub>O and 53°C, 98°C, 108°C and 116°C in D<sub>2</sub>O for He, Ne, Ar, Kr and Xe respectively [IAPWS]. The poor solubility of non-polar gases in water is due to positive free energy change attributed to the large negative entropy change caused by the structural enhancement of the water (ES) clusters; a conclusion reinforced by the enhanced heat capacity of these solutions. Counter-intuitively in spite of it forming stronger hydrogen bonds, D<sub>2</sub>O is a better solvent than H<sub>2</sub>O for non-polar gases, as it more easily forms the ES water clustering and therefore can accommodate the guest molecules more easily without breaking its hydrogen bonds. Additionally positively hydrating salts (e.g. LiCl) that destroy the water low-density ES clustering reduce the solubility ('salt out') of the non-polar gases whereas hydrophobic hydrating salts (e.g. tetramethylammonium chloride) that increase water low-density ES clustering stability also increase non-polar gas solubility ('salt in'). Some non-polar organic molecules also behave similarly to non-polar gases, but their increased size alters the clathrate structuring. Thus benzene has a solubility minimum, at a lower temperature than expected from above, at about 20°C [210],

Interestingly, the change in solubility of non-polar gases has a maximum (and their free energy of hydration has a minimum) with their diameter when it is about the same as that of the dodecahedral cavities (i.e. ~4 the icosahedral network [196]). The solubility behavior of larger hydrophobic molecules is discussed briefly elsewhere. It should also be noted that the solvent properties of liquid superheated water also change with temperature as its dielectric permittivity reduces towards that of common organic solvents as the temperature rises towards the critical point. [Back]

### **34 At low temperatures, the self-diffusion of water increases as the density and pressure increase.**

The increase in self-diffusion with density (within the range of about 0.9 g cm<sup>-3</sup> up to about 1.1 g cm<sup>-3</sup>) in contrast to normal liquids where increasing density decreases self-diffusion as the molecules restrict each other's movements. The density increase may be due to increasing temperature, below 4°C, at atmospheric pressure or due to increasing pressure at low temperatures. Liquids normally show reduced self-diffusion when they are squeezed but water at 0°C increases its diffusivity by 8% under a pressure of about 200 MPa [226] before any further increase in pressure reduces the diffusivity in common with the behavior of other liquids. The movement of water becomes restricted at low temperatures as the more open (lower density) structure produced on cooling (see above) is formed by stronger hydrogen bonding, which reduces the self-diffusion. The strength of the hydrogen bonding is a controlling influence in this anomalous region, where the hydrogen bond angles and the inter-molecular distances are strongly coupled and this order decreases on compression [169] due to the collapse of ES structures to CS structures. Simulation studies have shown that self-diffusion goes through a minimum as the density of water is reduced below about 0.9 g cm<sup>-3</sup> followed

an increase with further density reduction, as might be expected from most liquids [402], due to the disruption of the network as the now-stretched hydrogen bonds are broken. For the same reasons, the molecular rotational movement of water (reciprocal rotational relaxation time) also varies in direct proportion to the changes in self-diffusion (translational movement). [Back]

### **35 The thermal conductivity of water rises to a maximum at about 130°C and then falls.**

For most liquids the thermal conductivity (the rate at which energy is transferred down a temperature gradient falls with increasing temperature but this occurs only above about 130°C in liquid water [188]. As the temperature of water is lowered, the rate at which energy is transferred is reduced to an ever-increasing extent. Instead of the energy being transferred between molecules, it is stored in the hydrogen bonding fluctuations within the increasingly large clusters that occur at lower temperatures. Effectively the energy is used to shift the  $ES \rightleftharpoons CS$  equilibrium towards the CS structure, which possesses greater flexibility and has a greater number of bent hydrogen bonds, rather than the transference of kinetic energy. [Back]

### **36 Proton and hydroxide ion mobilities are anomalously fast in an electric field.**

The ionic mobilities of hydrogen ions and hydroxide ions at 362 and 206 ( $\text{nm s}^{-1}/(\text{V m}^{-1})$ ) at 25°C are very high compared with values for other small ions such as lithium ( $40(\text{nm s}^{-1})/(\text{V m}^{-1})$ ) and fluoride ( $57(\text{nm s}^{-1})/(\text{V m}^{-1})$ ). This is explained by the Grotthuss mechanism. [Back]

### **37 The heat of fusion of water with temperature exhibits a maximum at -17°C [15].**

This strange behavior has been determined from the variation in ice and water specific heat capacities ( $C_p$ ) is due to changes in the structuring of supercooled water. As the temperature is lowered from 0°C the hydrogen-bond strength of ice increases due to the reduction in their vibrational energy and this gives rise to an increasing difference (as temperature is lowered) between the enthalpy of the water and ice. At low temperatures (below about -17°C) the continued shift, with lowering temperature, in the supercooled water  $CS \rightleftharpoons ES$  equilibrium towards the ES structure reduces the enthalpy of the liquid water relative to the ice to the consequent increase in hydrogen-bond strength and this causes the drop in the heat of fusion with lowering temperature. [Back]

### **38 The dielectric constant is high (78.4 at 25°C).**

Polar molecules, where the centers of positive and negative charge are separated, possess dipole moment means that in an applied electric field, polar molecules tend to align themselves with the field. Although water is a polar molecule, its hydrogen-bonded network tends to oppose this alignment. The degree to which a substance does this is called its dielectric constant and, because water is exceptionally cohesive, it has a high dielectric constant. This allows it to act as a solvent for ionic compounds, where the attractive electric field between the oppositely charged ions is reduced by about 80-fold, allowing thermal motion to separate them into solution. On heating, the dielectric constant drops, and liquid water becomes far less polar, down to a value of about 6 at the critical point. The dielectric constant similarly reduces if the hydrogen bonding is broken by other means, such as strong electric fields.

Perhaps this property of water should not be considered anomalous as other small polar molecules form liquids also having high dielectric constants (e.g. the dielectric constant of liquid HCN is 158). [Back]

<sup>1</sup> Theoretical considerations concerning the ice nucleation site size gives estimates of 45,000 water molecules at -5°C down to 70 water molecules at -40°C [265]. Molecular dynamics studies show that these do not form a crystalline structure for crystallization to occur [347]. [Back]

<sup>2</sup> Although this behaviour is encountered with some other solute-solvent combinations (e.g. methane in *n*-heptane), the behavior is a more general property of water and deserves comment. [Back]

<sup>3</sup> There is evidence [157, 269, 274] that the first (clathrate) shell possesses stronger hydrogen bonding and weakens the hydrogen bonding out to the next shell. [Back]

<sup>4</sup> Pipes burst due to the rapid formation of a network of feathery dendritic ice enclosing water which then expands on freezing within a now restricted volume to generate the required pressure [354]. [Back]

Please submit any other explanations of water anomalies.

Water: Home | The anomalies of water | Water: Introduction | The icosahedral water clusters | School of Applied Science | S  
Bank University

This page was last updated by Martin Chaplin  
on **27 July, 2003**

▲ top ▲